# **Optimal Impedance Design for Piezoelectric Vibration Control**

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Abstract-Piezoelectric transducers are commonly used as strain actuators in the control of mechanical vibration. One control strategy, termed piezoelectric shunt damping, involves the connection of an electrical impedance to the terminals of a structurally bonded transducer. Many passive, non-linear, and semi-active impedance designs have been proposed that reduce structural vibration. This paper introduces a new technique for the design and implementation of piezoelectric shunt impedances. By considering the transducer voltage and charge as inputs and outputs, the design problem is reduced to a standard linear regulator problem enabling the application of standard synthesis techniques such as LQG,  $\mathcal{H}_2$ , and  $\mathcal{H}_\infty$ . The resulting impedance is extensible to multi-transducer systems, is unrestricted in structure, and is capable of minimizing an arbitrary performance objective. An experimental comparison to a resonant shunt circuit is carried out on a cantilevered beam. Previous problems such as ad-hoc tuning, limited performance, and sensitivity to variation in structural resonance frequencies are significantly alleviated.

#### I. INTRODUCTION

Active feedback control involves the use of sensors and actuators to minimize structural vibration. The vibration is sensed directly and used to derive an actuator voltage  $V_a$  counter-active to the applied disturbance. Typical vibration sensors include accelerometers, velocimeters, and strain sensors. The foremost difficulties associated with active feedback control are due mainly to the intrinsic nature of the plant G. Mechanical systems are of high order and contain a large number of lightly damped modes. The modeling and control design for such systems is well known to pose significant challenges. In addition, environmental variation of the structural resonance frequencies can further complicate the problem by compromising stability margins and restricting performance.

In active vibration control, and many other applications, piezoelectric transducers are used exclusively as either sensors or actuators. Dosch, Inman, Garcia [1] and Anderson, Hagood, Goodliffe [2] were able to demonstrate a technique now referred to as piezoelectric self-sensing, or sensori-actuation. By subtracting the capacitive voltage drop from the applied terminal voltage, a reconstruction of the internal piezoelectric strain voltage can be obtained. The reconstructed strain voltage can be employed as an active feedback sensor effectively eliminating the need for an auxiliary transducer. In addition to the usual problems associated with active feedback control, piezoelectric selfsensing systems are also highly sensitive to the transducer capacitance value. A sensing capacitance not perfectly

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matched to the transducer capacitance can result in significant errors in the strain estimation. If the estimate is used within a feedback control loop, such uncertainty may severely affect performance or cause instability. An attempt to address the problem of capacitance sensitivity can be found in [3], [4].

Another technique, first appearing in [5], termed shunt damping, involves the connection of an electrical impedance to the terminals of a piezoelectric transducer. Impedance designs have included resistors [6], inductive networks [7], [8], switched capacitors [9], switched networks [10], negative capacitors [11], and active impedances [12]. Shunt damping has a number of benefits and disadvantages when compared to active feedback control. Shunt circuits do not require a feedback sensor, and in some circumstances, may not require any support electronics or power supply at all. Typically, a shunt damping strategy involves a specific impedance structure which is designed to damp a number of targeted structural modes. Another advantage of shunt damping is that the circuits can be fine-tuned online to compensate for any modeling errors experienced during the design process. Automatic online tuning techniques have also been presented [13].

This paper presents a fully automatic technique for the design and implementation of piezoelectric shunt damping circuits. By viewing the transducer voltage and charge as inputs and outputs, the task of impedance design can be cast as a standard regulator problem. Synthesis techniques such as LQG,  $\mathcal{H}_2$ , and  $\mathcal{H}_\infty$  are readily applied to procure a suitable impedance. Unlike present methodologies, the impedance is unrestricted in structure, is multi-port for multi-transducer systems, and can be designed to meet any set performance specification within the flexibility of the synthesis process.

The following two sections, Impedance Synthesis, and Modeling, review the basic concepts of impedance synthesis and introduce a simple, charge based modeling technique for piezoelectric laminate structures. Section 4 outlines the control objectives and presents  $\mathcal{H}_2$ , and  $\mathcal{H}_\infty$  approaches to the task of impedance synthesis. Experimental results in Section 5 show superior performance to passive shunt damping circuits. The results and contributions are summarized in Section 6.

#### II. MODELING

With the aim of facilitating active shunt design, this section introduces a charge-based modeling technique for piezoelectric laminate structures.

Consider the piezoelectric laminate structure shown in Figure 1 (a). Through the use of a shunt patch driven by the



Fig. 1. A shunted multi-transducer structure (a). Synthetic implementation of the impedance (b).

voltage  $V_z$ , the goal is to suppress vibration resulting from two disturbances:  $V_a$ , the voltage applied to a disturbance patch, and f(r, t) a generally distributed external force. The implemented transfer function between the measured charge q and applied voltage  $V_z$  effectively presents an electrical impedance Z(s) to the transducer. The structure is disturbed by m transducers on the left side, and controlled by a further m collocated transducers on the other. Each piezoelectric transducer is modeled electrically as a capacitor  $C_{pm}$  in series with a strain-dependent voltage source  $v_{pm}$  [1], [6], [14]. The possibility for multiple transducers will be considered.

Expressions for the open-loop and shunted dynamics, can be found in [15]. The effect of a connected shunt impedance can be viewed as equivalent to a strain-feedback control system [16].

As discussed in [15], the system shown in Figure 1 (b) can be reduced to the input-output model shown in Figure 2. In conformance with the standard MIMO control formulation, the plant contains two sets of inputs: the disturbance signals w, and the control signals u. For the case under consideration, the disturbance and control signals are realized through a set of voltages  $V_a$  and  $V_z$  applied to a number of laminated piezoelectric patches. The system outputs  $V_p$ , d(r, t), and q, correspond respectively to the piezoelectric voltages induced in each shunt patch, the dynamic displacement measured at a point r, and the charge resident on each patch. The displacement signal d(r, t) is chosen as our performance variable z, while the measured charge q is our feedback variable y. Although the induced shunt piezoelectric voltages  $V_p$  are not required during the design, their inclusion aids in the modeling process. Given a specific s-impedance, the signal  $V_p$  also allows us to compute the equivalent collocated active feedback controller. A state-space realization is easily generated to



Fig. 2. The composite structural piezoelectric plant P.

represent the system P [15].

$$\dot{x} = \mathbf{A}x + \mathbf{B} \begin{bmatrix} V_a \\ V_z \end{bmatrix}$$
(1)  
$$\begin{pmatrix} V_p \\ d(r) \\ q \end{bmatrix} = \mathbf{C}x + \mathbf{D} \begin{bmatrix} V_a \\ V_z \end{bmatrix}$$

where

$$A = \begin{bmatrix} 0 & 1 & & & \\ -\omega_1^2 & -2\varsigma_i\omega_1^2 & & & \\ & & \ddots & & \\ & & & 0 & 1 \\ & & & -\omega_N^2 & -2\varsigma_N\omega_N^2 \end{bmatrix}$$
$$B = \begin{bmatrix} 0 & 0 & & & \\ F_1 & H_1 \\ \vdots & \vdots \\ 0 & 0 \\ F_N & H_N \end{bmatrix} C = \begin{bmatrix} E_1 & 0 & \cdots & E_N & 0 \\ 1 & 0 & \cdots & 1 & 0 \\ C_p E_1 & 0 & \cdots & C_p E_N & 0 \end{bmatrix}$$
$$D = \begin{bmatrix} D_{11} & D_{12} \\ D_{21} & D_{22} \\ D_{11}C_p & -C_p + D_{12}C_p \end{bmatrix}$$

where  $\varsigma_k$  are the damping ratios of each mode,  $\omega_k$  are the resonance frequencies,  $F_k$  and  $H_k$   $k \in \{1, 2, \dots N\}$  are the state-input weightings of each disturbance and shunt transducer. The vectors  $E_k$   $k \in \{1, 2, \dots N\}$  represent the contribution of each mode to the induced piezoelectric voltages.

As an alternative to the parameterized modeling approach presented above, a multi-variable time or frequency domain system identification technique could be employed to estimate the plant P directly from experimental data.

#### **III. S-IMPEDANCE CONTROL DESIGN**

Given the composite model discussed in Section II, the problem of designing an appropriate impedance can be cast as a standard  $\mathcal{H}_2$  or  $\mathcal{H}_\infty$  regulator problem. As shown in Figure 3, the regulator C(s) accepts the measured charge q to provide a control signal  $V_z$  counteractive to the applied disturbance  $V_a$ . The objective is to minimize the structural displacement d(r, t) subject to a weighting on the magnitude of the required terminal voltage  $V_z$ .

In an  $\mathcal{H}_2$  sense, the goal is to minimize the transfer function from an applied disturbance w to the performance



Fig. 3. The standard  $\mathcal{H}_2$  and  $\mathcal{H}_\infty$  design problem containing the composite plant P and a secondary performance signal weighting the applied shunt voltage  $V_z$ .

signal z, i.e. we seek to minimize

$$J = \left\| \frac{z(s)}{w(s)} \right\|_{2}$$

$$= \left\| \frac{d(r,s) + \mathbf{k}_{u} V_{z}(s)}{V_{a}(s)} \right\|_{2}.$$
(2)

where the  $\mathcal{H}_2$  norm  $||F(s)||_2$  of F(s) is defined as

$$\|F(s)\|_{2}^{2} = \frac{1}{2\pi} \int_{-\infty}^{\infty} tr \{F(j\omega)F(j\omega)'\} \ d\omega.$$
 (3)

By Parseval's equality, the optimal  $\mathcal{H}_2$  controller minimizes the expected root-mean-square (RMS) value of z. An optimal  $\mathcal{H}_2$  controller can be found through the solution of an algebraic Ricatti equation.

Disadvantages associated with  $\mathcal{H}_2$  and LQG methods include the unrealistic Gaussian disturbance model, and problems related to integral performance constraints.  $\mathcal{H}_{\infty}$ optimization and robust control, originally championed by Zames [17], is an alternative to  $\mathcal{H}_2$  and LQG methods.

Applying  $\mathcal{H}_{\infty}$  control to the problem of s-impedance synthesis involves finding a controller C(s) that minimizes

$$J = \left\| \frac{z(s)}{w(s)} \right\|_{\infty}$$

$$= \left\| \frac{d(r,s) + \mathbf{k}_u V_z(s)}{V_a(s)} \right\|_{\infty}.$$
(4)

where the  $\mathcal{H}_{\infty}$  norm  $||F(s)||_{\infty}$  of F(s) is defined as

$$\|F(s)\|_{\infty} = \max_{\omega} \ \bar{\sigma} \left(F(j\omega)\right) \tag{5}$$

where  $\bar{\sigma}$  denotes the maximum singular value.

In the time domain,  $\mathcal{H}_{\infty}$  control can be interpreted as minimizing the worst-case induced 2-norm of z, i.e.

$$\left\|\frac{z(s)}{w(s)}\right\|_{\infty} = \max_{w(t)\neq 0} \frac{\|z(t)\|_2}{\|w(t)\|_2}$$
(6)

where  $||f(t)||_{2}^{2} = \int_{0}^{\infty} \sum_{i} |f_{i}(t)|^{2} dt$ .

Closely resembling the solution to  $\mathcal{H}_2$  synthesis, an optimal  $\mathcal{H}_\infty$  controller can be found through the solution of an algebraic Ricatti equation. Linear Quadratic Gaussian methods (LQG) are also readily applied [15].

## IV. EXPERIMENTAL RESULTS

In the following sub-sections, an  $\mathcal{H}_{\infty}$  s-impedance controller is designed and applied experimentally to control a piezoelectric laminate cantilever beam.

# A. Experimental Apparatus

The experimental apparatus, shown in Figure 5 and pictured in Figure 4, consists of a uniform aluminium cantilever beam. Three piezoelectric transducers are laminated onto the front face and connected electrically in series to the voltage source  $V_z$ . A single collocated disturbance transducer, identical to each of the shunt transducers, is also mounted on the back face and driven with the disturbance voltage  $V_a$ . Details of the beam, piezoelectric transducers, and voltage amplifier can be found in [15]. The displacement measurement d(r,t) is acquired using a Polytec PSV300 scanning laser vibrometer.

### B. Parameter Identification

To determine the model parameters shown in equation (1), a simple optimization scheme is employed. From an initial guess,  $\omega_i$  and  $\zeta_i$ , are found through a simplex optimization based on the measured disturbance to displacement transfer function  $\frac{d(r,s)}{V_a(s)}$ , i.e.

$$\begin{bmatrix} \omega_k & \varsigma_k \end{bmatrix} = \arg \min \left\| \tilde{P}_{dV_a}(s) - P_{dV_a}(s) \right\|_2, \quad (7)$$

where  $P_{dV_a}(s)$  is the measured transfer function from an applied disturbance  $V_a(s)$  to the displacement d(r, s). With these parameters in hand, those remaining are determined from a final global optimization,

$$\arg \min \left\| \tilde{P}(s) - P(s) \right\|_{2, W} . \tag{8}$$

As gains from channel to channel vary greatly, a multivariable frequency weight W is required to normalize the cost of each error transfer function. After identification, a good correlation between the model and experimental data was observed [15].

In the following sections it will be of interest to examine the robustness of each control strategy to a change in the structural resonance frequencies. Experimentally, such variation is accomplished by affixing a mass 60 mm from the beam tip. The corresponding change in structural resonance frequencies is illustrated in Figure 6.

## C. Passive Shunt Design

For the sake of comparison, each LQG and  $\mathcal{H}_{\infty}$  shunt impedance will be judged against a traditional resonant piezoelectric shunt damping circuit applied to the same structure. A current-flowing shunt circuit [18] was designed and tuned to minimize the  $\mathcal{H}_2$  norm of the cantilever beam. The schematic and component values can be found in Figure 7 and Table I.

## D. $\mathcal{H}_{\infty}$ Shunt Design

As discussed in Section III, an  $\mathcal{H}_{\infty}$  s-impedance is designed to minimize the following cost function,

$$J = \left\| \frac{d(r,s) + k_u V_z(s)}{V_a(s)} \right\|_{\infty},\tag{9}$$

where  $k_u$ , the control signal weighting, was chosen to be  $3.2 \times 10^{-7}$ . A random auxiliary input of negligible influence was also included to avoid plant inversion. For a discussion on plant inversion and its avoidance, see Fleming 2004 [15].

The complex s-impedance of the resulting  $\mathcal{H}_{\infty}$  controller is plotted in Figure 8.

Examining the open- and closed-loop pole locations shown in Figure 9, the controller is clearly augmenting the system damping. Corresponding mitigation of the transfer function from an applied disturbance to the measured displacement can be seen in both the frequency domain, Figure 10, and the time domain, Figure 12. The magnitude of the first and second structural modes are reduced by 30.3 and 24.0 dB respectively. Damping ratios are increased from 0.00246 and 0.0011 to 0.0288 and 0.00766.

An unexpected feature of the s-impedance is its smooth frequency response; there are no localized peaks at the resonance frequencies. In contrast, active strain-, velocity-, or acceleration-feedback controllers characteristically apply a highly localized gain at the frequencies of structural resonance. In the advent of model variation, such localized behavior can result in considerable performance degradation. In order to examine system robustness, the nominal system is perturbed by adding a mass 60 mm from the beam tip. Aside from the disturbance to the underlying partial differential equation, the first and second resonance frequencies are decreased by 13.5 and 2.2 % respectively. The consequence on both passive and active shunt circuits is shown in Figure 11. While the  $\mathcal{H}_{\infty}$  shunt loses only 3.3 and 0.8 dB from its unperturbed attenuation of the first and second modes, the passive shunt loses 13.4 and 4.8 dB. Corresponding time domain results are shown in Figure 12.



Fig. 4. The cantilever beam.

#### V. CONCLUSIONS

A framework has been presented for the design of active shunt impedances. By viewing a piezoelectric laminate structure as a system with transducer voltage inputs and



Fig. 5. A front elevation of the cantilever beam. A single co-located disturbance transducer excited by the voltage  $V_a$ , is also mounted on the back face.



Fig. 6. The experimental frequency response (in decibels) from an applied disturbance voltage  $V_a$  (V) to the resulting tip displacement d (m). Free (- -), With Mass (—).

charge outputs, the task of shunt impedance design can be accomplished through the solution of a standard control problem e.g. by LQG,  $\mathcal{H}_2$ , or  $\mathcal{H}_\infty$  synthesis. The resulting controller, effectively the derivative of impedance, can be implemented directly with a voltage amplifier and charge measurement.

Although the fundamental goal in smart structure design is often to the augment system damping, this cannot be specified directly as an LQG,  $\mathcal{H}_2$ , or  $\mathcal{H}_\infty$  performance objective. The approach has been to achieve this indirectly through mitigation of the performance transfer function  $\frac{d(s)}{V_\alpha(s)}$ .

Experimentally, the active shunts have proven to introduce significant system damping, up to 30.3 dB attenuation of the first cantilever mode.

While achieving levels of performance previously only available through sensor-based feedback control, active shunt impedances are remarkably insensitive to variation in the structural resonance frequencies. A 13.5 % change in the first resonance frequency resulted in only a slight loss in performance. By comparison, the same variation had a disastrous consequence on the performance of a

- 1	10 101	$C_2$	$10 \ nr$
$L_1$	11690 H	$L_2$	348 H
$R_1$	$15 \ k\Omega$	$R_2$	$9 \ k\Omega$

TABLE I

COMPONENT VALUES OF THE CURRENT-FLOWING SHUNT CIRCUIT.



Fig. 7. A dual-mode current-flowing piezoelectric shunt damping circuit [18].

passive shunt damping circuit. Such sensitivity has limited the practical application of smart structures incorporating either active feedback or passive shunt vibration control systems.

Another well known problem associated with passive shunt damping is the lack of control influence. Given a lightly damped structure, even the small counteractive forces associated with passive shunt circuits can significantly increase system damping. Many practical mechanical structures naturally exhibit higher levels of damping. In such cases, passive piezoelectric shunt circuits are of limited use. As the amount of control influence associated with active shunt impedances is arbitrary, the possibility now exists for controlling more heavily damped systems. In such cases, the control voltage  $V_z$  is expected to become quite large. At high drive voltages it may be necessary to address the inherent piezoelectric hysteresis.

The reader will appreciate that the presented techniques are quite general and valid for structures incorporating multiple piezoelectric transducers. Although the application of sensor-based feedback control is well defined and feasible for structures with multiple sensors and actuators, the same can not be said for multi-transducer shunt circuits [16] . Present multi-transducer, multi-mode shunt circuits are simply a direct extension of single transducer shunt circuits. Each circuit is restricted to be independent and attached to a single transducer. If a single mode is to be targeted by two or more transducers, the task of tuning the shunt circuit can become extremely tedious. In addition to the complicated interaction between transducers at those frequencies, there are now as many more tuning parameters as there are transducers per mode. The design freedom afforded with active shunts not only eliminates the complicated task of tuning, but allows for full utilization of each patch. The resulting impedance is unstructured, multivariable, and able to exploit benefits that may arise from inter-transducer coupling.



Fig. 8. Complex s-impedance of the  $\mathcal{H}_\infty$  (—), and ideal negative capacitor (- -) shunt controller.



Fig. 9. The open- ((), and closed-loop (×) pole locations of the  $\mathcal{H}_{\infty}$  shunt controlled system.

Possible applications of active piezoelectric shunt impedances include sensor-less, high performance vibration control of acoustic panels, flexible structures, and positioning systems. Future work includes multi-transducer structures and restricted impedance design. The LQG and  $\mathcal{H}_{\infty}$  impedance designs contained negative reactive components and are unstable in a systems perspective. Although the connection of the transducer and control impedance is stable, an inherently stable controller is desirable. It is presently unclear if an unstable controller is necessary to result in effective vibration reduction.

#### VI. ACKNOWLEDGMENTS

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Fig. 10. The experimental (—), and simulated (- -),  $\mathcal{H}_{\infty}$  shunt controlled frequency responses (in decibels) from an applied disturbance voltage  $V_a$  (V) to the resulting tip displacement d (m). The open-loop response is also shown (—).



Fig. 11. The free (- -), and with-mass (—), passive (a) and  $\mathcal{H}_{\infty}$  shunt controlled (b) experimental frequency responses (in decibels) from an applied disturbance voltage  $V_a$  (V) to the resulting tip displacement d (m).

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Fig. 12. The free (left column) and with-mass (right column) tip displacement response d(m) to a step disturbance in  $V_a$ . Experimental open-loop (a), passive shunt controlled (b), and  $\mathcal{H}_{\infty}$  shunt controlled (c) systems.

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